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Crystal chemistry of Te(IV) fluorinated compounds

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Abstract

After a short review of the main results of our activities in the field of fluorinated compounds, a comparison of the crystal structures of the $M_x Te_v^{(IV)}F_z$ compounds is developed.

Considering only the shortest bonds (in the range 1.80–2.10 Å) the co-ordination polyhedra for Te atoms are quasi systematically TeF₅E square pyramids or more rarely TeF₄E disphenoids. With increasing M/Te ratio, the fluorinated network changes from bi-dimensional Te_{1.5}F_{6.5} layers of corner-sharing TeF₄E and TeF₅E polyhedra in RbTe₃F₁₃, to independent TeF₅E units in the MTeF₅ compounds, via finite Te₄F₁₈ groups of four basic units, in KTe₂F₉. Such an evolution emphasises the depolymerising character of the M atoms.

In all these structures the electronic lone pairs E of Te atoms are stereochemically active, pointing towards the empty spaces of the various networks: cavities, tunnels or inter-sheet spaces. © 2001 Elsevier Science B.V. All rights reserved.

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1. Introduction

Crystal structure determination of fluorinated materials by X-ray, electron or neutron diffraction experiments, has now become a quasi-routine scientific activity, correctly carried out in most of the concerned laboratories. On the contrary, the subtle analysis, the vivid description and consequently the full exploitation of the information collected are far from being as well performed, and deserve a more constant interest from our scientific community.

We have been involved for many years in this particularly fascinating field of solid state chemistry (whose predictive character is not the least interesting) by systematically comparing new structures solved with listed structural types, and by elaborating original defect models, in close relation with nonstoichiometry and with short- or long-range order.

In the case of fluoroindates and fluoroscandates $M_x In(Sc)F_{3+x}$ of monovalent metals we have shown, for the first time, that indium or scandium atoms, in addition to the classical octahedral co-ordination, can accept a sevenfold co-ordination (pentagonal bipyramid), and we have constructed a structural classification of these compounds [1–6].

In the case of fluorozirconates, more than 40 crystalline phases have been identified and structurally characterised [7–24]. Several homologous structural series have been described, as for example, the anion-excess fluorite-related $M_m X_{2m+2}$ and $M_m X_{2m+4}$ series [14], or the square antiprism layer structures [28]. An accurate analysis of the structure of α -ZrF₄ [29], β -BaZr₂F₁₀ [17] and β -BaZrF₆ [9], has allowed a significant breakthrough in the understanding of BaF₂–ZrF₄ glasses [17,25–27].

We have been strongly involved in the characterisation of structural defects in anion-excess fluorine-related phases like $M_{1-x}^{II}M_x'^{II+}F_{2+\alpha x}$ or $(M,M')(O,F)_{2+x}$ complex oxide-fluorides and $M(N,F)_{2+x}$ nitridefluorides.

In the nonstoichiometric disordered phases, the defects (interstitial anions, anion vacancies, relaxed anions) constitute finite clusters. We have shown that depending on the nature of the fluorite matrix and of the dopant cation, three kinds of clusters are observed: small 1:0:*n* clusters, cuboc-tahedral 8:12:0 or 8:12:1 clusters, and columnar 2n + 2:4n:2 or 2n + 2:3n:2 clusters. The stability of these clusters depends more on the relative size of dopant and host cations than on their charge [30–42]. These clusters interact in the highly doped phases so as to constitute more or less well ordered superstructures as for example the incommensurate or semi-commensurate vernier phases $M_m X_{2m+1}$ [43,44].

In mixed anion phases a O/F or N/F order is systematically observed: the smaller and lower charged anions are always located in the densiest part of the structure, and only a strict order can justify the existence of abnormally vernier phases such as Bi_3NF_6 [45] and $Lu_3O_5F_2$ [46].

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More recently, we have undertaken a systematic study in the solid state of the MF–TeF₄ (M: Li, Na, K, Rb, Cs, Tl, NH₄), MF₂–TeF₄ (M: Ba, Pb) and TeO₂–TeF₄ systems.

In the case of the TeO₂–TeF₄ system [47], two original oxidefluorides TeOF₂ [48] and Te₂O₃F₂ [49] have been synthesised. Both compounds exhibit strong structural relationships with respectively the α - and β -TeO₂ polymorphs. In both structures, bond valence calculations indicate a perfect O/F order: O atoms logically occupy the equatorial positions in the anionic polyhedra and are bridging atoms; F atoms, on the contrary, occupy axial positions and are non-bridging.

In the MF_2 -TeF₄ systems, in addition to a fluorite-type solid solution, which in the case of M: Pb is an excellent anionic conductor, two stoichiometric phases, with compositions $MTeF_6$ and MTe_2F_{10} , have been found and studied structurally.

In the MF–TeF₄ systems, three kinds of stoichiometric compounds have been isolated: MTeF₅ (M: Li, Na, K, Rb, Cs, Tl, NH₄), MTe₂F₉ (M: K), MTe₃F₁₃ (M: Rb) and perhaps MTe₄F₁₇ (M: Cs) [50]. It is interesting to note that, in the TeO₂–TeF₄ system, as well as in three MF–TeF₄ systems (M: K, Rb, Cs), very large glassy domains (about 55–95 mol% TeF₄) have been observed, demonstrating the strong glass-forming character of TeF₄. The structural study of such glasses, particularly the oxidefluorides, which exhibit glass transition at temperatures lower than 0°C, would be of fundamental interest.

2. The $M_x Te_v F_z$ crystal structures

In all the known structures the Te atoms are fivefold coordinated (square pyramid) approximately at the centre of a distorted octahedron whose one open corner is occupied by its stereochemically active electronic lone pair E (Fig. 1a). Such a basic building unit, which is also present in the TeF₄ crystal structure [51], and whose geometry does not notably change with the nature of the M cation or of the network

Fig. 1. Spatial view of (a) the TeF₅E polyhedron (KTe₂F₉, MTeF₅ and RbTe₃F₁₃) (b) the TeF₄₊₂E polyhedron (KTe₂F₉ and RbTe₃F₁₃).

formed seems, therefore, to be the most striking feature in the crystal chemistry of Te(IV) in fluorinated media.

In the KTe₂F₉ [52] and RbTe₃F₁₃ [50] structures, respectively, one half and two thirds of the Te atoms are fourfold co-ordinated (distorted disphenoid) approximately at the centre of a distorted trigonal bipyramid whose one equatorial corner is occupied by the lone pair E (Fig. 1b).

All the Te–F distances are short, in the range 1.81–2.08 Å. Because of the stereochemical activity of the lone pair E:

- in the TeF₅E octahedra, the Te-F_{ax} bond is systematically shorter than the four Te-F_{eq} ones, and the F-Te-F angles are always less than 90°;
- in the TeF₄E trigonal bipyramids, the two Te-F_{eq} bonds are shorter than the two Te-F_{ax} ones, and the F_{ax}-Te-F_{ax}, F_{eq}-Te-F_{eq} and F_{ax}-Te-F_{eq} angles are always less than the ideal respective values: 180, 120 and 90°.

In the case of the TeF₄E polyhedra observed in the KTe₂F₉ [52] and RbTe₃F₁₃ [50] crystal structures, if we take into account the two extra anions at distances 2.52–2.65 Å, the co-ordination polyhedra can then be considered as very distorted pentagonal bipyramids TeF₆E, whose one equatorial corner is occupied by the lone pair E (Fig. 1b).

Let us see now how the spatial distribution and the linking of these basic structural units develop when

- for a given stoichiometry the nature of the M cation changes;
- the Te/M ratio increases.

2.1. Te/M = 1

This corresponds to the MTeF₅ [53–58] and MTeF₆ [59,60] crystal structures. In all these structures the TeF₅E basic units are independent.

As shown in a forthcoming paper [61], all the MTeF₅ structures are made up of a mix in equal number of independent M^+ cations and $(TeF_5E)^-$ anions, and they are all topologically related to the CsCl type, i.e. the TeF₅E units are located within M_8 more or less distorted cubic cavities. Such a cavity is emphasised, in Fig. 2, which represents the KTeF₅ crystal structure (orthorhombic symmetry, space group Pbcm [54]). It is worth noting that in each case:

- due to the strong stereochemical activity of the lone pairs E, the TeF₅E units are set off-centre toward an edge of the M₈ cavities.
- the distortion of the two primitive cubic arrays is all the more important since the geometrical fit between the (TeF₅E)⁻ anions and the M⁺ cations is bad.

In the PbTeF₆ structure [59], which can be related to the ReO₃ type, the three-dimensional network consists of double columns of edge- and corner-sharing PbF₇E polyhedra, connected one to each other by sharing corners with independent TeF₅E units.





Fig. 2. Projection along $[0 \ 0 \ 1]$ of the KTeF₅ crystal structure showing the K₈ cubic cavities. Arrows visualise the direction toward which point the lone pairs E (in the upper level).

In the BaTeF₆ structure [60], which is related to the tysonite (LaF₃) type, the Ba atoms and the TeF₅E units are distributed at the corners of a distorted fcc array. By sharing triangular faces the BaF₁₁ polyhedra constitute infinite columns connected by rows of independent TeF₅E units.

2.2. Te/M = 2

This corresponds to the KTe_2F_9 and $BaTe_2F_{10}$ crystal structures.

The BaTe₂F₁₀ structure (space group C2/c, a = 17.311 Å, b = 4.832 Å, c = 10.233 Å, $\beta = 111.60^{\circ}$, Z = 6) [62] is very simple and corresponds to a slightly distorted cubic close packing of Ba atoms and independent TeF₅E units. As clearly shown in Fig. 3a, it can be described as the

regular and alternate stacking along $[1\ 0\ 0]$ of 4⁴ plane nets of Ba atoms and layers of TeF₅E polyhedra, according to the sequence ...Ba–TeF₅E–TeF₅E–Ba…. In fact, it corresponds to the regular repetition along $[1\ 0\ 0]$ of BaTe₂F₁₀E₂ three-layer-wide sheets, deriving from the KOsF₆ structure (rhombohedrally distorted CsCl type), each sheet being shifted $1/2\langle 1\ 1\ 0\rangle$ CsCl one from the other.

If we only take into account the shortest Te–F distances, the KTe₂F₉ crystal structure (space group P-1, a = 4.757 Å, b = 9.333 Å, c = 9.467 Å, $\alpha = 108.82^{\circ}$, $\beta = 95.79^{\circ}$, $\gamma = 96.11^{\circ}$, Z = 2) [52], corresponds to the same sequence of $4^2 \times 3^3$ plane nets of K atoms and of layers of TeF₄E and TeF₅E independent units (Fig. 3b). Such an analogy justifies the close geometrical relationships which exist between the two unit cells: $a(BaTe_2F_{10}) \approx 2c(KTe_2F_9)$, $b(BaTe_2F_{10}) \approx a(KTe_2F_9)$, $c(BaTe_2F_{10}) \approx b(KTe_2F_9)$ and



Fig. 3. Comparison of the crystal structures of (a) $BaTe_2F_{10}$; (b) KTe_2F_9 . Arrows visualise the direction toward which point the lone pairs E.

 $V(\text{BaTe}_2\text{F}_{10}) \approx V(\text{KTe}_2\text{F}_9)$. The doubling of *a*, and, therefore, of the unit cell volume in $\text{BaTe}_2\text{F}_{10}$ results from the periodic shift ($\tau = 1/2\langle 110\rangle$ CsCl, see Fig. 3a) of the KOsF₆-like sheets, not observed in the KTe₂F₉ structure. Despite this analogy, the two structures are different. Whereas in the $\text{BaTe}_2\text{F}_{10}$ structure the TeF₅E units are independent, in the KTe₂F₉ one, the TeF₄E and TeF₅E units are connected one to another via F(6) and F(9) anions (see Fig. 3b) producing discrete Te₄F₁₈ units.

2.3. Te/M = 3

This case is limited to the $RbTe_3F_{13}$ crystal structure [50].

As shown in Fig. 4, in this structure the Rb atoms and independent TeF₅E and TeF₄E units are distributed at the corners of a face centred orthorhombic array. That it corresponds to the regular stacking along [0 0 1] of alternate square 4⁴ plane nets of Rb atoms and of layers of independent TeF₅E and TeF₄E units, according to the sequence \cdots Rb–Te(2)F₄E–Te(1)F₅E–Te(2)F₄E–Rb \cdots . If we take into account the two weak bonds (Te(2)F = 2.65 Å), the structure can, therefore, be described as a succession along [0 0 1] of 4⁴ square layers of Rb atoms and of two independent Te_{1.5}F_{6.5} sheets of corner-sharing Te(1)F₅E and Te(2)F₄E units, separated by the lone pairs E.



Fig. 4. Projection onto $[0\ 1\ 0]$ of the RbTe₃F₁₃ crystal structure. Broken lines indicate the weak Te(2)–F bonds, and arrows visualise the direction toward which point the lone pairs E.

3. Conclusions

In all the fluorinated compounds, as in all the corresponding complex oxides or oxidefluorides, the lone pairs E of Te(IV) atoms are stereochemically active and fill the empty spaces of the various networks: cavities, tunnels, inter-sheet spaces.

With increasing M/Te ratio, the fluorinated network changes from infinite $Te_{1.5}F_{6.5}$ layers of corner-sharing TeF₅E and TeF₄E polyhedra in RbTe₃F₁₃ to independent TeF₅E units in the MTeF₅ structures, via finite Te₄F₁₈ groups of four basic units in KTe₂F₉. Such an evolution clearly



Fig. 5. Projection of the TeF_4 crystal structure down [1 0 0], showing the chains of corner-sharing TeF_5E polyhedra.

emphasises the depolymerising role of the M atoms. A similar evolution should be observed in the corresponding fluorinated glasses. Moreover, it is worth pointing out that as soon as M/Te > 1, the two constituents KF and TeF₄ tend to segregate structurally in alternate sheets, the thickness of the TeF_x sheets increasing with the M/Te ratio. From this point of view, the determination of the CsTe₄F₁₇ crystal structure would be of great interest. It is also interesting to note the great analogies of the infinite chains of TeF₅E and TeF₄E polyhedra which, by crossing constitute the Te₃F₁₃ sheets (Fig. 4), with the zigzag chains TeF₄ present in the TeF₄ structure [51] (Fig. 5).

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